

**MATH 4033 • Spring 2018 • Calculus on Manifolds**  
**Problem Set #3 • Tensors and Differential Forms • Due Date: 08/04/2018, 11:59PM**

1. Show that the Lie derivative of a 1-form

$$\alpha = \sum_i \alpha_i du^i$$

along a vector field  $X = \sum_j X^j \frac{\partial}{\partial u_j}$  is locally expressed as:

$$\mathcal{L}_X \alpha = \sum_{i,j} \left( X^i \frac{\partial \alpha_j}{\partial u_i} + \alpha_i \frac{\partial X^i}{\partial u_j} \right) du^j.$$

2. Consider a smooth manifold  $M^n$  and the following symmetric  $(2,0)$ -tensors on  $M$ :

$$g = \sum_{i,j} g_{ij} du^i \otimes du^j \qquad h = \sum_{i,j} h_{ij} du^i \otimes du^j$$

where  $(u_1, \dots, u_n)$  are local coordinates of  $M$ . Suppose the matrix  $[g_{ij}]$  is positive definite (all its eigenvalues are positive).

- (a) Let  $(v_1, \dots, v_n)$  be another local coordinates system, and the local expression of  $g$  in terms of this system be:

$$g = \sum_{i,j} \tilde{g}_{ij} dv^i \otimes dv^j.$$

Express  $\tilde{g}_{ij}$  in terms of  $g_{ij}$ , and then show that the matrix  $[\tilde{g}_{ij}]$  is also symmetric and positive definite.

- (b) Show that the quantity  $\frac{\det[h_{ij}]}{\det[g_{ij}]}$  is independent of local coordinates, i.e.

$$\frac{\det[h_{ij}]}{\det[g_{ij}]} = \frac{\det[\tilde{h}_{ij}]}{\det[\tilde{g}_{ij}]}.$$

3. Define a 2-form  $\omega$  on  $\mathbb{R}^3$  by

$$\omega = x dy \wedge dz + y dz \wedge dx + z dx \wedge dy.$$

- (a) Express  $\omega$  in spherical coordinates  $(\rho, \theta, \varphi)$  defined by:

$$\begin{aligned} x &= \rho \sin \varphi \cos \theta \\ y &= \rho \sin \varphi \sin \theta \\ z &= \rho \cos \varphi \end{aligned}$$

- (b) Compute  $d\omega$  in both rectangular and spherical coordinates.  
(c) Let  $S^2$  be the unit sphere  $x^2 + y^2 + z^2 = 1$  in  $\mathbb{R}^3$ . Consider the inclusion map  $\iota : S^2 \rightarrow \mathbb{R}^3$ . Compute the pull-back  $\iota^* \omega$ , and express it in terms of spherical coordinates.  
(d) Let  $\delta = dx \otimes dx + dy \otimes dy + dz \otimes dz$  be a symmetric  $(2,0)$ -tensor on  $\mathbb{R}^3$ . Compute  $\iota^* \delta$ , and express it in terms of spherical coordinates.

4. The purpose of this exercise is to show that any closed 1-form  $\omega$  on  $\mathbb{R}^3$  must be exact. Let

$$\omega = P(x, y, z) dx + Q(x, y, z) dy + R(x, y, z) dz$$

be a closed 1-form on  $\mathbb{R}^3$ . Define  $f : \mathbb{R}^3 \rightarrow \mathbb{R}$  by:

$$f(x, y, z) = \int_{t=0}^{t=1} (xP(tx, ty, tz) + yQ(tx, ty, tz) + zR(tx, ty, tz)) dt.$$

- (a) Show that  $df = \omega$ .  
 (b) Point out exactly where you have used the fact that  $d\omega = 0$  in your solution to (a).  
 (c) Explain why your solution to (a) would fail if the domain of  $\omega$  is  $\mathbb{R}^3 \setminus \{p\}$  (where  $p$  is a fixed point in  $\mathbb{R}^3$ ).
5. Consider  $\mathbb{R}^4$  with coordinates  $(t, x, y, z)$ , which is also denoted as  $(x_0, x_1, x_2, x_3)$  in this problem. Denote  $*$  to be the Minkowski Hodge-star operator on  $\mathbb{R}^4$  (see P.104).
- (a) Compute each of the following:

$$\begin{array}{lll} *(dt \wedge dx) & *(dt \wedge dy) & *(dt \wedge dz) \\ *(dx \wedge dy) & *(dy \wedge dz) & *(dz \wedge dx) \end{array}$$

- (b) The four Maxwell's equations are a set of partial differential equations that form the foundation of electromagnetism. Denote the components of the electric field  $E$ , magnetic field  $B$ , and current density  $J$  by

$$\begin{aligned} E &= E_x i + E_y j + E_z k \\ B &= B_x i + B_y j + B_z k \\ J &= j_x i + j_y j + j_z k \end{aligned}$$

All components of  $E$ ,  $B$  and  $J$  are considered to be time-dependent. Denote  $\rho$  to be the charge density. The four Maxwell's equations assert that:

$$\begin{array}{ll} \nabla \cdot E = \rho & \nabla \cdot B = 0 \\ \nabla \times E = -\frac{\partial B}{\partial t} & \nabla \times B = J + \frac{\partial E}{\partial t} \end{array}$$

We are going to convert the Maxwell's equations using the language of differential forms. We define the following analogue of  $E$ ,  $B$ ,  $J$  and  $\rho$  using differential forms:

$$\begin{aligned} E &= E_x dx + E_y dy + E_z dz \\ B &= B_x dy \wedge dz + B_y dz \wedge dx + B_z dx \wedge dy \\ J &= (j_x dy \wedge dz + j_y dz \wedge dx + j_z dx \wedge dy) \wedge dt + \rho dx \wedge dy \wedge dz \end{aligned}$$

Define the 2-form  $F$  by  $F := B + E \wedge dt$ . Show that the four Maxwell's equations can be rewritten in an elegant way as:

$$dF = 0 \qquad d(*F) = J.$$